

Time Reversal Violating Resonance Energy Shifts Revisited

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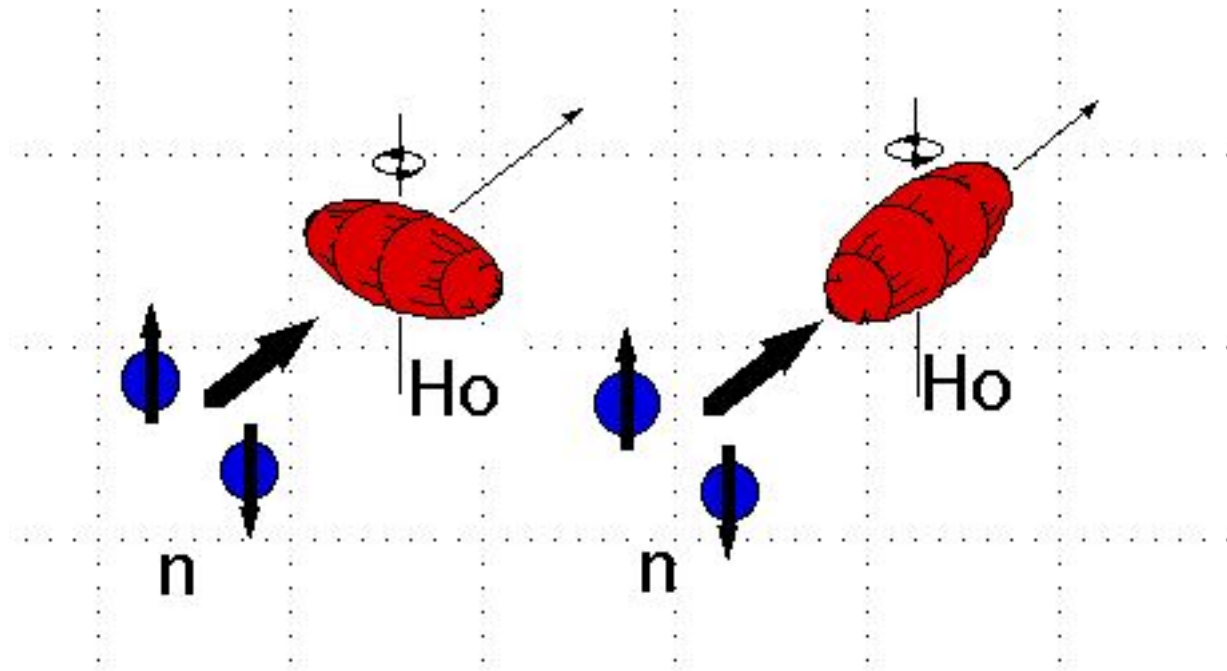
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June 16, 2003

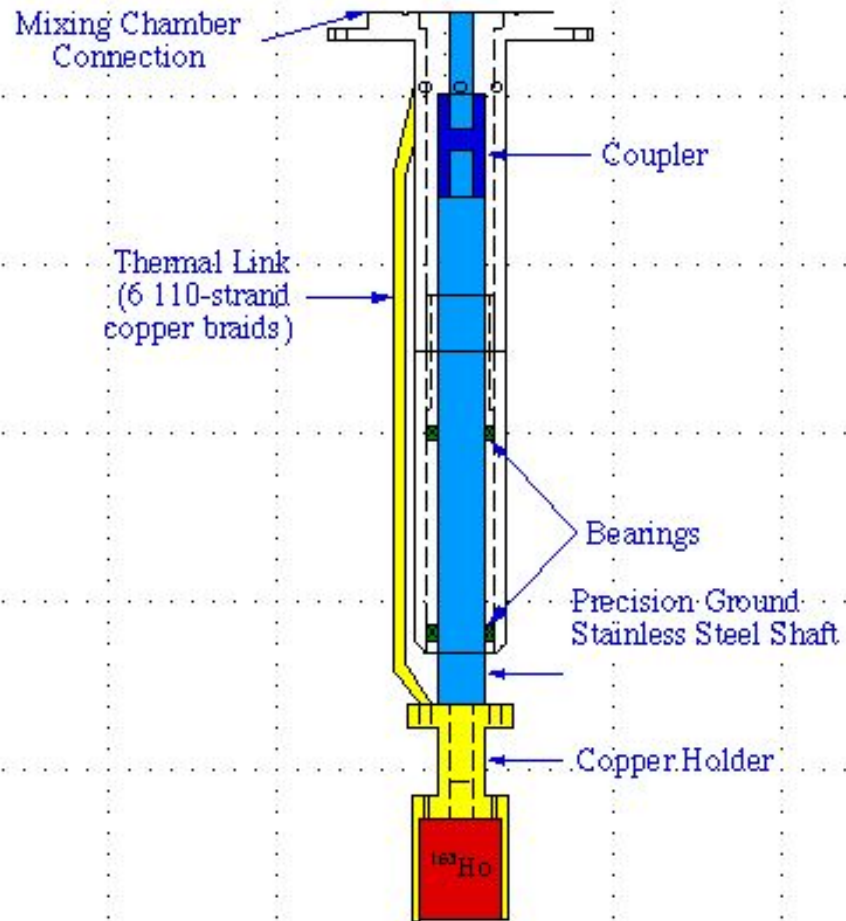
Outline

- *Polarized neutron transmission tests of TRV*
- *TRV and resonance energy spacings –early work*
- *Polarized neutron capture energy shifts – another way?*
- *T even “fake” energy shifts*
- *Statistical issues and ultimate bounds*

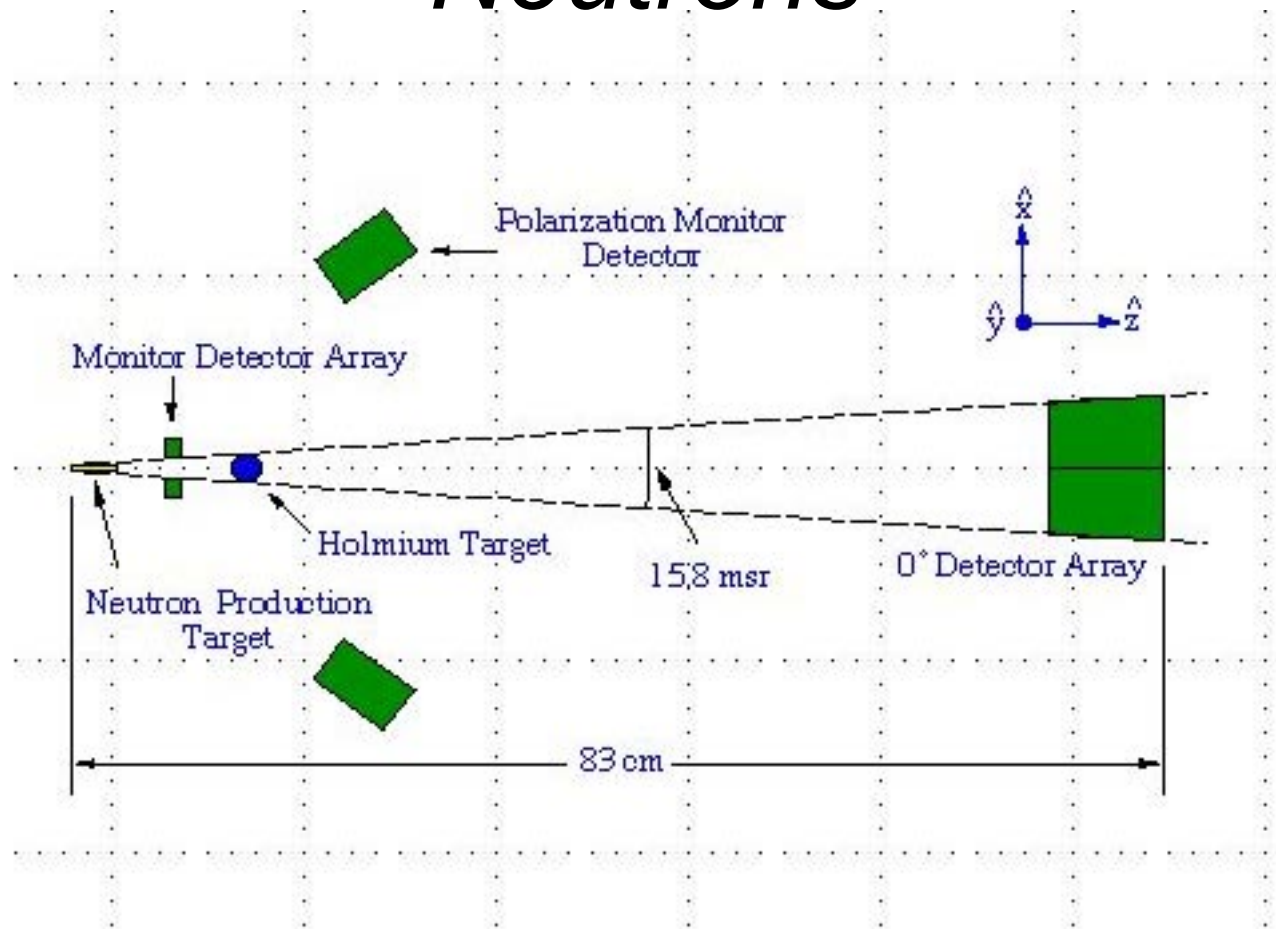
*FC Test with polarized neutrons
and aligned holmium nuclei*



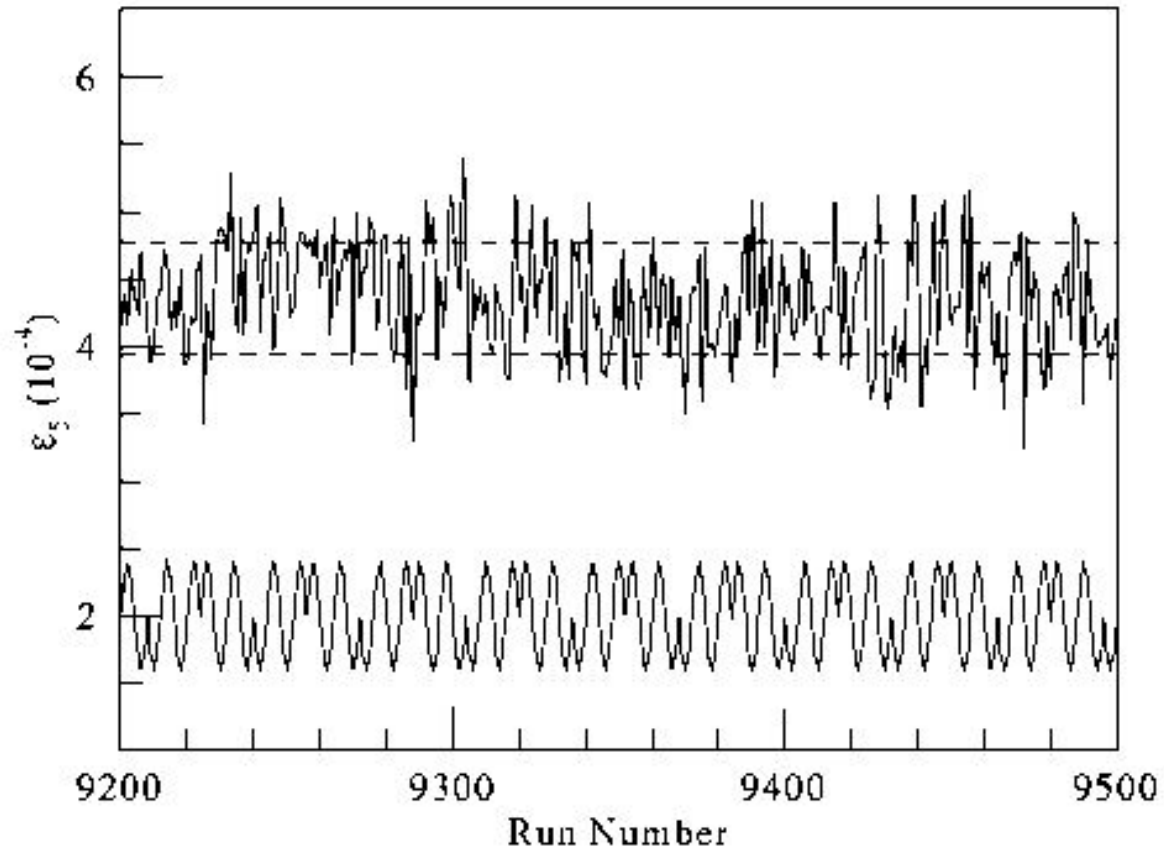
Mounting Apparatus



TUNL FC Test with 6 MeV Neutrons



P-even TRV transmission asymmetry



Results for MeV neutron tests

TC (89) 7,11 MeV $\varepsilon \leq 10^{-3}$

FC (91,97) 6 MeV $\varepsilon = 1.1 \pm 1.0 \times 10^{-6}$ (400 \times TDB)

1/A suppression $\rightarrow \check{g}_\rho = 2.3 \pm 2.1 \times 10^{-2}$

edm's rule, but for a P-even TRV scenario where they set no bounds, see Kurylov et al (PR D63, 2001)

COSY plans 200 MeV test in p-d scattering (2004)

Why no resonance tests?

- *TC – sequential interactions mimic the signal: strong rotation by F_M followed by weak attenuation by F_P –target problem*
- *FC – sequential interactions not a problem: need a beam line and a holmium target!*
- *“for whatever reason this field has been dead in the water for some time” (R. Golub on edm’s)*

Is there another way?

- *T-odd energy shifts in neutron capture correlations*
- *Do not need a polarized target!*
- *P-odd with polarized neutrons (Flambaum and Sushkov 85)*
- *P-even with unpolarized neutrons (Barabanov 93)*
- *But (caution) not a null test*

Early ideas on P-even TRV from energy level spacings

French et al (87) shell model study of GOE versus GUE:

$H = H + i U, \quad = |\langle v_t \rangle / \langle v \rangle|$ ratio of residual interaction 2 body matrix elements

=0 : GOE (no T-violation)

=1 : GUE (maximal T-violation)

From analysis of Th energy levels, they conclude

$$\leq 2 \times 10^{-3} \text{ (99\% CL)} \Rightarrow \check{g}_\rho \leq 2 \times 10^{-1}$$

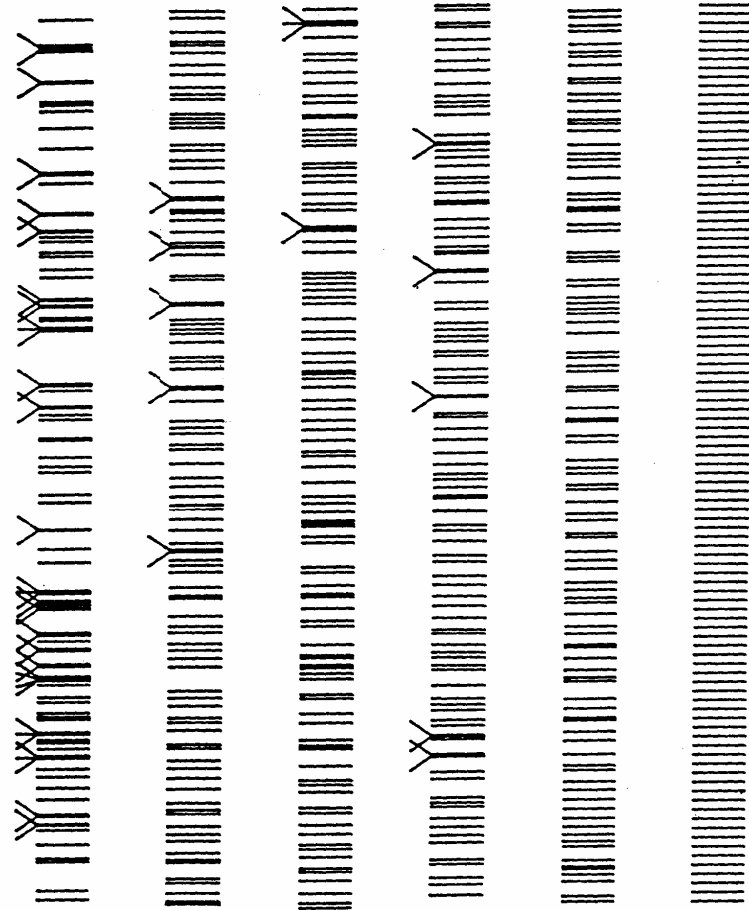
where $\check{g}_\rho \approx$ Haxton et al 94), is dimensionless ratio of T-odd to T-even ρ coupling constants

Match the spectra contest!

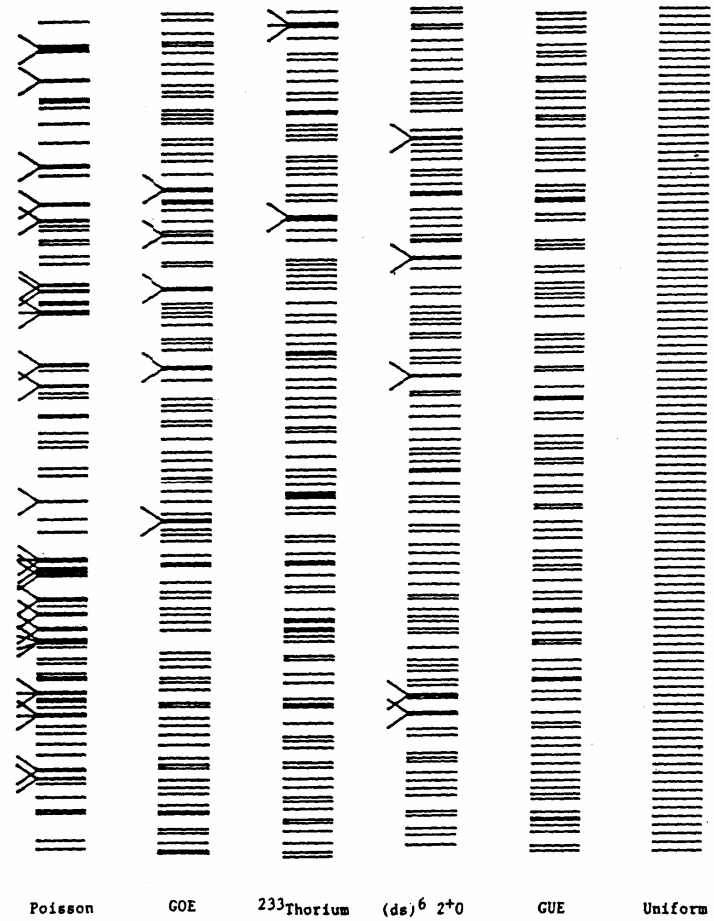
- Six spectra* – only one is “real”

1. Poisson
2. Uniform
3. GOE
4. GUE -
5. Shell model study of (s-d) $6 2^+$ levels
6. Thorium

*Thanks to A. Pandey



Only #3 is real



P-even TRV energy shift study from Dubna

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Testing T -Odd, P -Even Interactions with γ Rays from Neutron p -Wave Resonances

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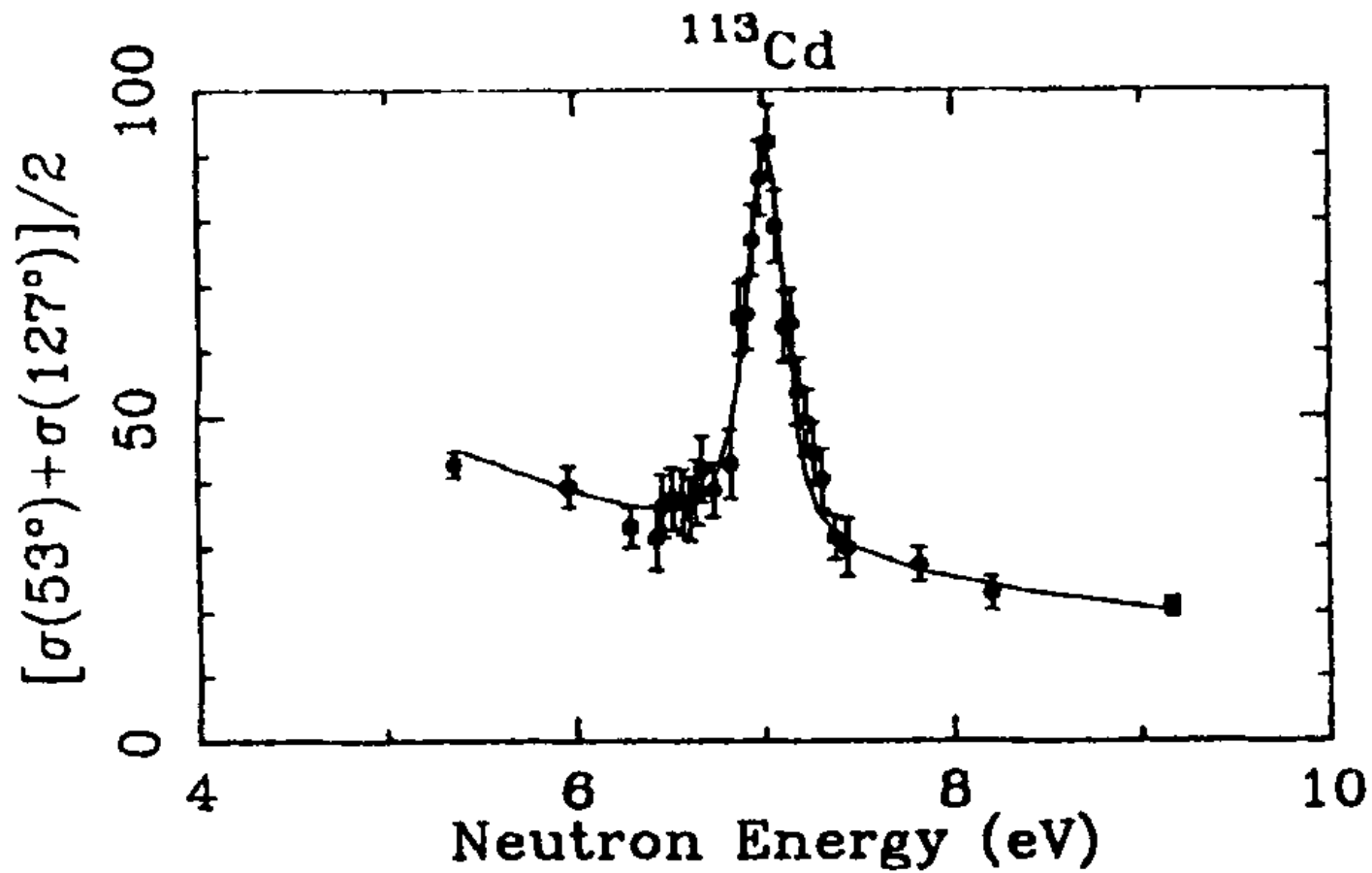
Los Alamos National Laboratory, Los Alamos, New Mexico 87545

(Received 13 October 1992)

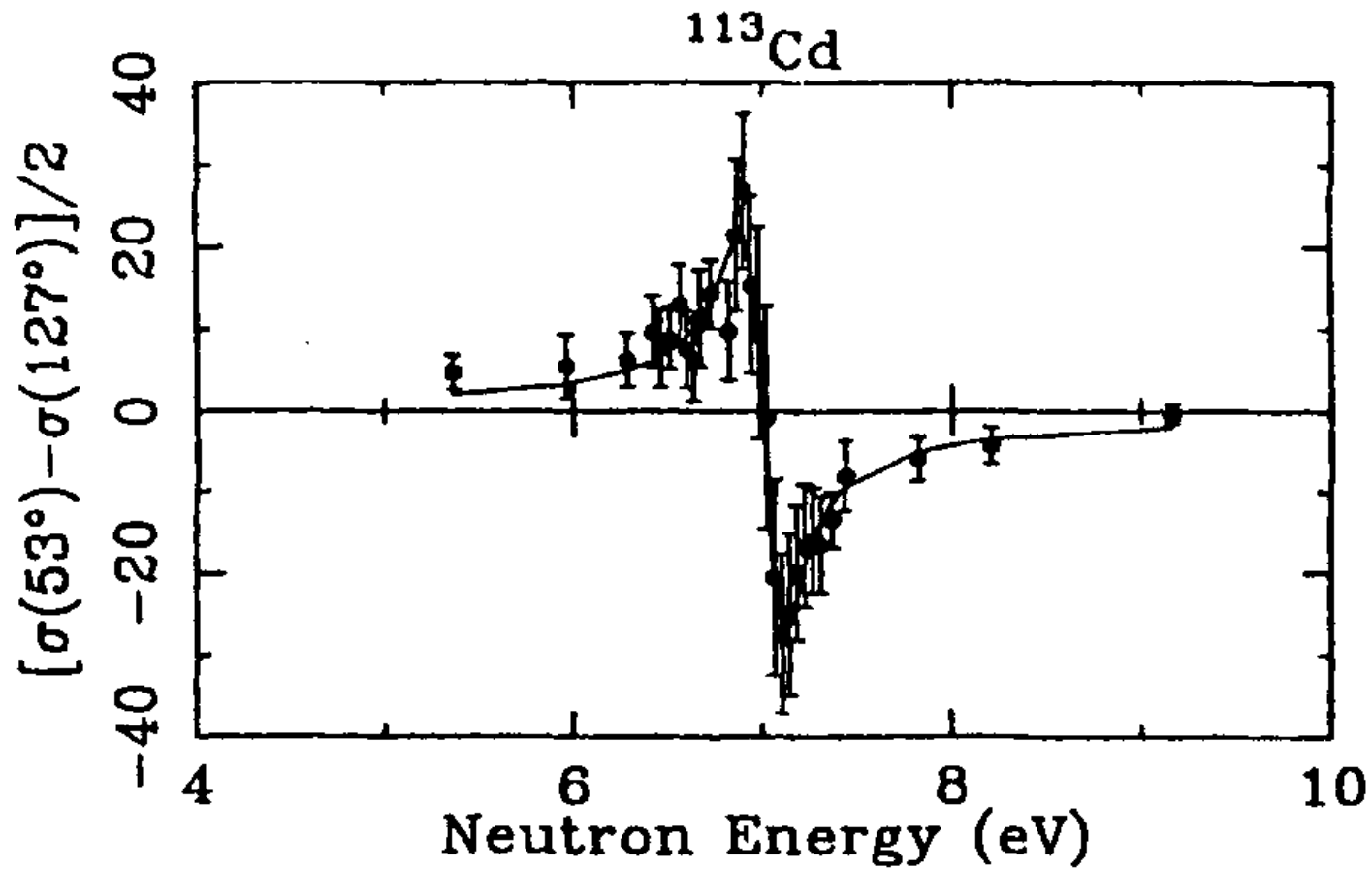
A new method for the study of time reversal violation is described. It consists of measurements of the forward-backward asymmetry in individual gamma-ray transitions resulting from unpolarized neutron capture in p -wave resonances. An experiment with a ^{113}Cd target performed at the Dubna pulsed neutron source has been analyzed and a limit on the time reversal odd, parity even interaction extracted. The possibility of experiments using the powerful pulsed neutron source at Los Alamos is considered.

PACS numbers: 25.40.Ny, 11.30.Er

Resonance energy from the yield curve



Zero crossing of the fore-aft asymmetry



Results in ^{113}Cd

- $E = 7.0412 \pm 0.0073 \text{ eV}$
- $\Delta E = -0.0016 \pm 0.0062 \text{ eV}$
- Conclusion is $\alpha < 10^{-4}$ (good precision)
- But single case, so difficult to draw a statistically meaningful bound

Energy shift signature of P-odd TRV

F and S proposal NP A435(84) p 363

*Seventeen angular and polarization correlations (8 P-even
7 P-odd) in (n,) reactions near threshold (s and p waves,
E1 and M1 amplitudes), coefficients a_i .*

*Focus on $s_n \mathbf{k}$ ray asymmetry with a transversely
polarized neutron beam (cf NPDG apparatus) at a p-wave
resonance showing large P violation*

P-odd TRV energy shift

Given the ratio of T-odd to T-even matrix elements

$$= |\langle s | v_{pt} | p \rangle / \langle s | v_p | p \rangle|$$

Then TRV shifts the zero crossing of the correlation term from $E=E_p$ by an amount proportional to

$$a_9 \sim \text{Re } 1/(E - E_p + i \Gamma/2) + \text{Im } 1/((E - E_p + i \Gamma/2))$$

Also true for circular polarization measured with an unpolarized beam

FS conjecture is surprising- is it true?

Revisit this issue by considering in order:

Two resonance model

Complex phases for neutron and γ -ray

Role of distant states

Statistical issues

Experimental issues

In FS notation, focus on $a_9' = a_9 - a_{12}/3$

Main conclusion: result is correct, but what are limitations?

Assume equal s and p wave amplitudes at $E=E_p$ (a requirement for large parity mixing)

Definitions

$$= (E - E_p) / (\hbar^2 / 2m)$$

$$= |\langle s | v_{pt} | p \rangle / \langle s | v_p | p \rangle|$$

Findings so far

Correction term $b \sim O(1/D)$ even in two resonance approximation: $a_0' \sim a + b$

Corrections due to complex phases in s channel are negligible $O(1/D)/\sqrt{N}$, N number of open channels

Corrections due to mixing of distant s -wave states are $O(1/D)$

Corrections due to mixing of distant p -wave states are suppressed by extra factor $O(1/D)$ over s wave mixing

*Probability distribution for γ is a convolution of a Gaussian and a Cauchy distribution – Voigt profile-
Bayesian analysis underway (EDD). Need multiple cases!*

Experimental Issues

- *NPDG apparatus exists – ideal for this measurement (measure transverse to avoid fore-aft P-even asymmetries)*
- *Ultimate sensitivity bounded by Γ/D – of order 10^{-3} ?*
- *Best experimental precisions quoted at Dubna and LANSCE are of order 6 meV – Q: is 1 meV possible?*

Conclusions

Energy shifts ARE a signature of P-odd TRV; a TRV neutron resonance experiment without a polarized target is feasible

However it's not a null test; correction terms (cf FSI) exist to $O(\lambda/D) \sim 10^{-3}$

Bring on the SNS?

The Fundamental Neutron Physics Beamline at the Spallation Neutron Source

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Contest: Identify the level diagrams

- Six spectra* – only one is “real”
 1. Poisson
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CP nonconservation or T (time reversal) noninvariance was observed in the decay of neutral kaons long ago [1]. Attempts to search for this phenomenon in other processes have failed. Recently some possible tests using neutron optics experiments were discussed in a paper by Weidenmuller [2] and references therein. Interest in the neutron-nucleus interaction was stimulated by the discovery [3,4] of large P - (parity-) nonconserving effects near neutron p -wave resonances. These effects appeared to be enhanced by 3 to 6 orders of magnitude compared to the single-particle estimate $Gm_n^2 \sim 10^{-7}$ (in units of $\hbar = c = 1$ and $G = 10^{-5} m_n^{-2}$). The enhancement was explained in terms of the P -nonconserving mixing of compound nuclear states with opposite parity (mixing of the p -wave resonance with the nearest s -wave resonances [5,6]). The concept of the spreading width of the weak interaction was brought into practice based on recent experimental results [7].

It was shown by Bunakov and Gudkov [8,9] that possible T -noninvariant effects may be enhanced near p -wave resonances in the same manner as P -nonconserving effects. It has been suggested that a search be made for P -conserving, T -noninvariant effects in polarized neutron transmission through aligned nuclear targets [10-12]. Such an effect would occur due to an $s: [k \times I](k \cdot I)$ term in the neutron-nucleus elastic forward scattering amplitude. Here s and I are the neutron and nucleus spins and k is neutron momentum. Bunakov [9] showed that the maximum value of the experimental polarization asymmetry in the cross section β is connected to the T -noninvariant matrix element v_p^T and to the p -wave resonance level spacing D_p by the approximate relation

$$\beta \approx v_p^T / D_p. \quad (1)$$

Here $iv_p^T = (HT)_{12} = -(HT)_{21}$ is the purely imaginary matrix element of the T -noninvariant, P -conserving interaction between compound nuclear wave functions of the considered p -wave resonance (labeled 1) and its nearest-neighbor p -wave resonance (labeled 2). Estimates given by Bunakov [9] indicate the enhancement

factor may be as big as 10^3 . The estimate is that $v_p^T / D_p \sim 10^3 \phi$, where ϕ is roughly the strength of the T -noninvariant, P -conserving nuclear interaction relative to the T -invariant one. The new theoretical limits on ϕ are very low [13]. The experimental upper limit on ϕ is between 10^{-3} and 10^{-4} , e.g., as given in Ref. [9] or in the recent aligned target experiment [14]. Therefore, in the framework of Bunakov's hypothesis, any experimental measurement of β near a p -wave resonance with accuracy better than 10^{-1} may decrease this limit.

It has been noted [15] that the value of v_p^T / D_p in p -wave resonances may be obtained from a quite different kind of experiment, namely, from the measurement of the forward-backward asymmetry in the yield of gamma rays from individual transitions in unpolarized neutron capture reactions. Such measurements have already been made [16-18] for p -wave resonances in targets of ^{113}Cd and ^{117}Sn . Here we present the results of our analysis of the ^{113}Cd capture experiment after giving the details of our approach to the analysis which is based on the hypothesis of Ref. [9].

In our approach to the calculation of experimental observables we will use an S -matrix scattering formalism. We wish to start from the expression

$$\delta S_J = S_J(1 \frac{1}{2} \rightarrow 1 \frac{1}{2}) - S_J(1 \frac{1}{2} \rightarrow 1 \frac{1}{2}),$$

where $S_J(j, j')$ is the scattering matrix element which corresponds to the transition of p -wave neutrons ($l = l' = 1$) in a resonance with spin J from a channel with neutron total angular momentum $j = \frac{1}{2}$ to a channel with $j' = \frac{3}{2}$ and vice versa. Using the explicit forms for the scattering matrix elements of Ref. [19] we get

$$\delta S_J = \frac{2 \text{Im}[g_n(1 \frac{1}{2}) g_n^*(1 \frac{3}{2})]}{E - E_p + i\Gamma_p/2}. \quad (2)$$

Here E_p and Γ_p are the energy and the total width of the p -wave resonance and $g_n(l, j)$ is the neutron partial width amplitude $[\Gamma_n(l, j) = |g_n(l, j)|^2]$. By definition

$$g_n = \alpha \langle \varphi_{JM} | \Psi_{JM} \rangle, \quad (3)$$